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Anti-synchronization of Duffing Double-Well Chaotic Oscillators via Integral Sliding Mode Control

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Abstract: Chaos theory has a manifold variety of applications in science and engineering. There are many systems in nature with several stable states, which are separated by energy barriers. When the system can move along the stable states, its dynamics can become quite complex. A simple mechanical model that depicts some of these complex dynamical features is the famous Duffing double-well oscillator (1918). This paper gives a summary description of the Duffing double-well chaotic oscillator. Next, new control results are obtained for the global chaos anti-synchronization of the identical Duffing double-well chaotic oscillators via integral sliding mode control (ISMC). MATLAB plots have been shown to illustrate the phase portraits of the Duffing double-well chaotic oscillator and the global chaos anti-synchronization of Duffing double-well chaotic oscillators via integral sliding mode control.

Keywords: Chaos, chaotic systems, chaos control, anti-synchronization, Duffing oscillator, mechanical system, oscillators, stable states, nonlinear oscillations, sliding mode control.

1. Introduction

A dynamical system is called *chaotic* if it satisfies the three properties: boundedness, infinite recurrence and sensitive dependence on initial conditions [1-4].

Chaos theory has a lot of applications in science and engineering [5]. Chaos theory has applications in dynamo systems [6-12], memristors [13-16], nonlinear oscillators [17-30], Tokamak systems [31-32], finance system [33], cellular neural networks [34-39], chemical reactors [40-50], neurology [51-58], population biology systems [59-67], etc.

A simple mechanical model that depicts some of these complex dynamical features is the famous Duffing double-well oscillator ([68], 1918). This paper gives a summary description of the Duffing double-well chaotic oscillator. Next, new results are obtained for the global chaos synchronization of the identical Duffing double-well chaotic oscillators via integral sliding mode control (ISMC). Sliding mode control is a popular control technique used in the control and synchronization of chaotic systems [69-73]. MATLAB plots have been shown to illustrate the phase portraits of the Duffing double-well chaotic oscillators via integral sliding mode control.

2. Duffing double-well chaotic oscillator

Duffing double-well chaotic oscillator [68] is described by the 2-D dynamics

$$\begin{cases} \dot{x} = y \\ \dot{y} = x - x^3 - ay + F\cos(\omega t) \end{cases}$$
(1)

The system (1) is *chaotic* when the system parameters are chosen as

$$F = 0.7, \ a = 0.5, \ \omega = 1$$
 (2)

For numerical simulations, we take the initial conditions

$$x(0) = \Theta.5, \quad y(0) \quad -0.5 \tag{3}$$

The 2-D phase portrait of the Duffing double-well chaotic oscillator (1) is depicted in Figure 1.



Figure 1. The 2-D phase portrait of the Duffing double-well chaotic oscillator

3. Anti-synchronization of the identical Duffing double-well chaotic oscillators

In this section, we use integral sliding mode control (ISMC) to achieve global chaos anti-synchronization of the identical novel Duffing double-well chaotic oscillators. We use Lyapunov stability theory to prove the main result derived in this section for the global chaos synchronization of the Duffing double-well chaotic oscillators.

As the master system, we consider the Duffing double-well chaotic oscillator given by

$$\begin{cases} \dot{x}_{1} = y_{1} \\ \dot{y}_{1} = x_{1} - x_{1}^{3} - ay_{1} + F\cos(\omega t) \end{cases}$$
(4)

where x_1, y_1 are the states and a, F, ω are constant positive parameters.

As the slave system, we consider the Duffing double-well chaotic oscillator given by

$$\begin{cases} \dot{x}_2 = y_2 + u_x \\ \dot{y}_2 = x_2 - x_2^3 - ay_2 + F\cos(\omega t) + u_y \end{cases}$$
(5)

where x_2, y_2 are the states and u_x, u_y are integral sliding mode controls to be determined.

The anti-synchronization error between the systems (4) and (5) is defined by

$$\begin{cases} e_x(t) = x_2(t) + x_1(t) \\ e_y(t) = y_2(t) + y_1(t) \end{cases}$$
(6)

The error dynamics is obtained as

$$\begin{cases} \dot{e}_{x} = e_{y} + u_{x} \\ \dot{e}_{y} = e_{x} - ae_{y} - x_{2}^{3} - x_{1}^{3} + 2F\cos(\omega t) + u_{y} \end{cases}$$
(7)

Based on the sliding mode control theory [74], the integral sliding surface of each error variable is defined as follows:

$$\begin{cases} s_x = \left[\frac{d}{dt} + \lambda_x\right] = \left[\int_0^t e_x(\tau)d\tau\right] & e_x + \lambda_x \int_0^t e_x(\tau)d\tau \\ s_y = \left[\frac{d}{dt} + \lambda_y\right] = \left[\int_0^t e_y(\tau)d\tau\right] & e_y + \lambda_y \int_0^t e_y(\tau)d\tau \end{cases}$$
(8)

The derivative of each equation in (8) yields

$$\begin{cases} \dot{s}_x = \dot{e}_x + \lambda_x e_x \\ \dot{s}_y = \dot{e}_y + \lambda_y e_y \end{cases}$$
(9)

The Hurwitz condition is satisfied if λ_x and λ_y are positive constants.

Based on the exponential reaching law [74], we set

$$\begin{cases} \dot{s}_x = -\eta_x \operatorname{sgn}(s_x) - k_x s_x \\ \dot{s}_y = -\eta_y \operatorname{sgn}(s_y) - k_y s_y \end{cases}$$
(10)

Comparing equations (9) and (10), we get

$$\begin{cases} \dot{e}_x + \lambda_x e_x & -\eta_x = \operatorname{sgn}(s_x) - k_x s_x \\ \dot{e}_y + \lambda_y e_y & -\eta_y = \operatorname{sgn}(s_y) - k_y s_y \end{cases}$$
(11)

Using Eq. (7), we can rewrite Eq. (11) as follows:

$$\begin{cases} e_y + u_x + \lambda_x e_x & -\eta_x = \operatorname{sgn}(s_x) - k_x s_x \\ e_x - a e_y - x_2^3 - x_1^3 + 2F \cos(\omega t) + u_y + \lambda_y e_y & -\eta_y = \operatorname{sgn}(s_y) - k_y s_y \end{cases}$$
(12)

From Eq. (12), the control laws are obtained as follows:

$$\begin{cases} u_x = -e_y - \lambda_x e_x - \eta_x \operatorname{sgn}(s_x) - k_x s_x \\ u_y = -e_x + ae_y + x_2^3 + x_1^3 - 2F \cos(\omega t) - \lambda_y e_y - \eta_y \operatorname{sgn}(s_y) - k_y s_y \end{cases}$$
(13)

Theorem 1. The Duffing double-well chaotic oscillators (4) and (5) are globally and asymptotically anti-synchronized for all initial conditions by the integral sliding mode controller (13), where the constants $\lambda_x, \lambda_y, \eta_x, \eta_y, k_x, k_y$ are all positive.

Proof. This result is proved using Lyapunov stability theory [75].

We consider the following quadratic Lyapunov function

$$V(s_x, s_y) = \frac{1}{2} \left(s_x^2 + s_y^2 \right)$$
(14)

where s_x , s_y are as defined in (8).

The time-derivative of V is obtained as

$$V = s_x \dot{s}_x + s_y \dot{s}_y \tag{15}$$

Substituting from Eq. (10) into (15), we get

$$\dot{V} = s_x \left[-\eta_x \operatorname{sgn}(s_x) - k_x s_x\right] + s_y \left[-\eta_y \operatorname{sgn}(s_y) - k_y s_y\right]$$
(16)

Simplifying Eq. (16), we obtain

$$\dot{V} = -\eta_x |s_x| - k_x s_x^2 - \eta_y |s_y| - k_y s_y^2$$
(17)

Since $k_x, k_y > 0$ and $\eta_x, \eta_y > 0$, it follows from (17) that \dot{V} is a negative definite function.

Thus, by Lyapunov stability theory [75], it follows that $(s_x, s_y) \rightarrow (0, 0)$ as $t \rightarrow \infty$.

Hence, it is immediate that $(e_x, e_y) \rightarrow (0, 0)$ as $t \rightarrow \infty$. This completes the proof.

4. Numerical Simulations

We use classical fourth-order Runge-Kutta method in MATLAB with step-size $h = 10^{-8}$ for solving the system of differential equations (4) and (5), when the integral sliding mode controller (13) is implemented.

The parameter values of the Duffing double-well oscillators are taken as in the chaotic case, viz.

$$a = 0.5, \ \ \omega \quad 1, \ \ E \quad 0.7$$
 (18)

We take the sliding constants as

$$\lambda_x = \lambda_y = 0.1 = \eta_x = \eta_y = 0.1 = k_x \quad k_y \quad 30 \tag{19}$$

We take the initial conditions of the Duffing double-well chaotic oscillator (4) as

$$x_1(0) = 3.7, y_1(0) = 12.8$$
 (20)

We take the initial conditions of the Duffing double-well chaotic oscillator (5) as

$$x_2(0) = 4.1, y_2(0) = 16.4$$
 (21)

Figures 2-3 show the anti-synchronization of the Duffing double-well chaotic oscillators (4) and (5).

Figure 4 shows the time-history of the anti-synchronization errors e_x, e_y .



Figure 2. Anti-synchronization of the states x_1 and x_2



Figure 3. Anti-synchronization of the states y_1 and y_2



Figure 4. Time-history of the anti-synchronization errors e_x, e_y

5. Conclusions

In this paper, we first gave a summary description of the Duffing double-well chaotic oscillator. Next, new results were obtained for the global chaos anti-synchronization of the identical Duffing double-well chaotic oscillators via integral sliding mode control (ISMC). MATLAB plots were shown to illustrate all the main results derived in this research work for the Duffing double-well chaotic oscillator.

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